

RESEARCH ARTICLE

Homogenization of a Pseudoparabolic System

Małgorzata Peszyńska, Ralph Showalter* and Son-Young Yi

Department of Mathematics, Oregon State University, Corvallis, OR 97331

(v3.3 released May 2008)

Pseudoparabolic equations in periodic media are homogenized to obtain upscaled limits by asymptotic expansions and two-scale convergence. The limit is characterized and convergence is established in various linear cases for both the classical binary medium model and the highly-heterogeneous case. The limit of vanishing time-delay parameter in either medium is included. The double-porosity limit of Richards' equation with dynamic capillary pressure is obtained.

Keywords: homogenization, pseudoparabolic equations, fractured porous media, dynamic capillary pressure

AMS Subject Classification: Primary 35B27, 35K70; Secondary 74Q10, 76S05.

1. Introduction

Pseudoparabolic equations arise in a range of applications from radiation with time delay [1], degenerate double-diffusion and heat-conduction models [2, 3] and resolution of ill-posed problems [4] through recently developed applications in level set methods [5] and models of lightning propagation [6]. They were first analyzed in [7–9]; see [10] for an extensive review and bibliography. Here we are interested in a degenerate pseudoparabolic equation arising from modeling dynamic capillary pressure in unsaturated flow; specifically, we study the case of flow in heterogeneous media in which the coefficients are periodic on a fine scale.

The classical Richards equation for flow through a partially-saturated porous medium with porosity $\phi(x)$ and permeability $K(x)$ takes the form

$$\phi(x) \frac{\partial u(t, x)}{\partial t} + \nabla \cdot K(x) \frac{k_w(u(t, x))}{\mu_w} \nabla (P_c(u(t, x)) - \rho G D(x)) = 0, \quad (1)$$

where u denotes saturation, and gravitational effects depend on depth $D(x)$ and (constant) density ρ . Here $k_w(u)$, $P_c(u)$ denote relative permeability and capillary pressure relationships, respectively. This standard model follows from Darcy's law extended to multiphase flow and conservation of mass [11, 12] with the assumption that atmospheric pressure of air is constant. The model has been analyzed in [13–15] and elsewhere.

The experimental determination of the pressure-saturation relationship $p = -P_c(u)$ is based on the assumption that this is an instantaneous process, although

*Corresponding author. Email: show@math.oregonstate.edu

in reality it requires substantial time to approach an equilibrium before measurements can be taken. This led to the introduction of *dynamic capillary pressure* [16] in which $P_c(u)$ is replaced by $P_{c,dyn}(u) \equiv P_c(u) - \tau \frac{\partial u}{\partial t}$ with $\tau > 0$. Other dynamic models had been introduced earlier [17, 18]; see [19–22] for supporting experimental evidence. A similar model was derived in [23] by homogenization from standard two-phase models with special interface conditions.

The dynamic capillary pressure model of [16] leads to the nonlinear pseudoparabolic equation

$$\begin{aligned} \phi(x) \frac{\partial u(t, x)}{\partial t} + \nabla \cdot K(x) \frac{k_w(u(t, x))}{\mu_w} \nabla (P_c(u(t, x)) - \rho G D(x)) \\ - \nabla \cdot K(x) \frac{k_w(u(t, x))}{\mu_w} \nabla \tau(x) \frac{\partial u(t, x)}{\partial t} = 0. \end{aligned} \quad (2)$$

When written in terms of pressure $u \mapsto -P_c(u)$ (see Section 4) and linearized about a known solution u_0 , with $\kappa(x) \equiv K(x) \frac{k_w(u_0)}{\mu_w}$, ϕ replaced by $\phi \frac{\partial u}{\partial p} |_{u_0}$ and τ by $\frac{\tau}{\phi}$, the equation (2) takes the form

$$\phi(x) \frac{\partial u(t, x)}{\partial t} - \nabla \cdot \kappa(x) \nabla (u(t, x) + \tau(x) \phi(x) \frac{\partial u(t, x)}{\partial t}) = \nabla \cdot \kappa(x) \rho G D(x). \quad (3)$$

If the convective term is dropped, *i.e.*, set $D(x) = 0$, we obtain

$$\phi(x) \frac{\partial u(t, x)}{\partial t} - \nabla \cdot \kappa(x) \nabla (u(t, x) + \tau(x) \phi(x) \frac{\partial u(t, x)}{\partial t}) = 0. \quad (4)$$

In realistic porous media there is substantial variation of $\phi(x)$ and $K(x)$, as well as the nonlinear relationships $k_w(\cdot)$, $P_c(\cdot)$, $\tau(\cdot)$ in (2). Consequently the coefficients in linearized models (3) and (4) vary similarly. In this paper we derive homogenized models for (2) and (4), and in particular for the special case of *binary media* in which $\phi(x)$, $K(x)$, $\tau(x)$ and consequently $\kappa(x)$ oscillate between two respective constant values. See [24, 25] for further discussion of heterogeneous dynamic capillary pressure models, references and numerical results.

The multiscale analysis is aided by the structure of the *pseudoparabolic system*

$$\phi(x) \frac{\partial u(t, x)}{\partial t} + \frac{1}{\tau(x)} (u(t, x) - v(t, x)) = 0, \quad (5a)$$

$$- \nabla \cdot (\kappa(x) \nabla v(t, x)) + \frac{1}{\tau(x)} (v(t, x) - u(t, x)) = 0, \quad x \in \Omega. \quad (5b)$$

This system is equivalent to a single equation: if we eliminate v we obtain the *pseudoparabolic equation* (4) for the variable $u(t, x)$; v satisfies a similar equation. It is supplemented with corresponding boundary and initial conditions. Here we take homogeneous Dirichlet boundary conditions

$$v(t, s) = 0, \quad \text{a.e. } s \in \partial\Omega, \quad (5c)$$

and the initial condition

$$\phi(x) u(0, x) = \phi(x) u_*(x), \quad \text{a.e. } x \in \Omega. \quad (5d)$$

The well-posedness of the system (5) follows from very general assumptions on the coefficients and initial function. The following suffices for our purposes here.

Theorem 1.1: *Assume that functions $\phi(\cdot)$, $\kappa(\cdot)$, $\tau(\cdot) \in L^\infty(\Omega)$ are given, each with a strictly positive lower bound, and let $u_*(\cdot) \in L^2(\Omega)$. Then there is a unique pair $u(\cdot) \in H^1((0, T); L^2(\Omega))$ and $v(\cdot) \in L^2((0, T); H_0^1(\Omega))$ such that $u(0, \cdot) = u_*(\cdot)$ and*

$$\int_{\Omega} \left(\phi(x) \frac{\partial u(t, x)}{\partial t} \varphi(x) + \frac{1}{\tau(x)} (u(t, x) - v(t, x)) (\varphi(x) - \psi(x)) + \kappa(x) \nabla v(t, x) \cdot \nabla \psi(x) \right) dx = 0 \quad (6)$$

for all $\varphi(\cdot) \in L^2(\Omega)$ and $\psi(\cdot) \in H_0^1(\Omega)$.

Corresponding results hold under much more general conditions of non-negativity of the coefficients. See [10, 26–29]. The initial value u_* need be chosen only with $\phi(\cdot)^{1/2} u_*(\cdot) \in L^2(\Omega)$. Also, the *a-priori* estimates show explicitly that $u - v \rightarrow 0$ as $\tau \rightarrow 0$.

Our objective is to homogenize the system (5) and thereby the corresponding pseudoparabolic equation (4) when the coefficients depend (periodically) on a small parameter ε . The precise description of coefficients will follow below. Bensoussan, Lions and Papanicolaou [30] briefly investigated the homogenization of pseudoparabolic equations as an example for which the limiting problem is of a different type, and perhaps *non-local*, not even a partial differential equation. (See Chapter II, Section 3.9, pp. 318, 338.) We shall see below that this occurs when certain variables are eliminated or *hidden*. The limited regularity and estimates for solutions of the corresponding pseudoparabolic equation (4) makes the homogenization more delicate. Only in special cases is there a purely upscaled limit.

In Section 2 we obtain the formal asymptotic expansion of the solution for the linear equation (4) in the classical case and find the dependence of the limit on ϕ and τ . The analysis and homogenization of the linear system (5) by two-scale convergence is developed in Section 3 for ε -periodic binary coefficients and includes cases of $\tau \rightarrow 0$ with parabolic or first-order kinetic systems as limits. Finally, Section 4 contains the asymptotic expansion for a nonlinear highly-heterogeneous case arising from Richards' equation with dynamic capillary pressure.

2. Asymptotic Expansion

First we introduce periodic coefficients into the pseudoparabolic system (5) and use formal asymptotic expansions to obtain the limiting problem as the period scale $\varepsilon > 0$ tends to zero. Let Y denote the unit cube in \mathbb{R}^N , let there be given the Y -periodic functions $\phi(y)$, $\tau(y)$, $\kappa(y)$ and then define $\phi^\varepsilon(x) = \phi(\frac{x}{\varepsilon})$, $\tau^\varepsilon(x) = \tau(\frac{x}{\varepsilon})$, $\kappa^\varepsilon(x) = \kappa(\frac{x}{\varepsilon})$. The three functions ϕ^ε , τ^ε , κ^ε are the respective ε -periodic coefficients in (5), so the corresponding solution u^ε , v^ε to (5) depends on ε . We write these as formal asymptotic expansions

$$u^\varepsilon(t, x) = \sum_{p=0}^{\infty} \varepsilon^p u_p(t, x, y), \quad v^\varepsilon(t, x) = \sum_{p=0}^{\infty} \varepsilon^p v_p(t, x, y), \quad y = \frac{x}{\varepsilon}, \quad (7)$$

with each $u_p(t, x, \cdot)$, $v_p(t, x, \cdot)$ being Y -periodic.

Substitute (7) into (5) and collect terms by powers ε^p for $p \geq -2$. Note that

the gradient $\nabla = \nabla_x + \frac{1}{\varepsilon}\nabla_y$ is used in calculations where $y = x/\varepsilon$. The ordinary differential equation (5a) gives (at $p = 0$)

$$\phi(y) \frac{\partial u_0(t, x, y)}{\partial t} + \frac{1}{\tau(y)}(u_0(t, x, y) - v_0(t, x, y)) = 0.$$

The initial condition will always be assumed to be independent of the local variable, $y \in Y$.

The procedure for the elliptic equation (5b) is standard [30–32]. Equating to zero the coefficient of ε^{-2} in the expansion of (5b) gives

$$-\nabla_y \cdot \kappa(y) \nabla_y v_0(t, x, y) = 0, \quad y \in Y.$$

With the Y -periodic boundary conditions on v_0 , we conclude that $\nabla_y v_0(t, x, y) = 0$, and so $v_0 = v_0(t, x)$ is independent of $y \in Y$. From the combined coefficients of ε^{-1} in the expansion of (5b) we obtain

$$-\nabla_y \cdot \kappa(y) (\nabla_y v_1(t, x, y) + \nabla_x v_0(t, x)) - \nabla_x \cdot \kappa(y) \nabla_y v_0(t, x) = 0.$$

The last term is null, so the function $v_1(t, x, y)$ is the solution of an elliptic periodic boundary-value problem on Y , and we can represent it in terms of Y -periodic solutions $\omega_j(y)$ of the *cell problem* (see (17))

$$-\nabla_y \cdot \kappa(y) (\nabla_y \omega_j + e_j) = 0, \quad j = 1 \dots N.$$

This representation $v_1(t, x, y) = \sum_{j=1}^N \omega_j(y) \frac{\partial}{\partial x_j} v_0(t, x)$ (up to a function of x) will be used to compute the effective tensor κ^* below. Finally, collecting terms with ε^0 in the expansion of (5b) gives

$$\begin{aligned} & -\nabla_y \cdot \kappa(y) (\nabla_y v_2 + \nabla_x v_1) \\ & -\nabla_x \cdot \kappa(y) (\nabla_x v_0(t, x) + \nabla_y v_1(t, x, y)) + \frac{1}{\tau(y)} (v_0(t, x) - u_0(t, x, y)) = 0. \end{aligned}$$

Integrate this equation over Y . The first term vanishes due to Y -periodicity of each v_r , and the second becomes the effective elliptic contribution with the tensor κ^* . The third term gets averaged to yield the second equation of the system

$$\phi(y) \frac{\partial u_0(t, x, y)}{\partial t} + \frac{1}{\tau(y)} (u_0(t, x, y) - v_0(t, x)) = 0, \quad (8a)$$

$$-\nabla \cdot \kappa^* \nabla v_0(t, x) + \int_Y \frac{1}{\tau(y)} (v_0(t, x) - u_0(t, x, y)) dy = 0, \quad (8b)$$

the first being copied from above. The effective tensor κ^* is obtained in this calculation as $\kappa_{ij}^* = \int_Y \kappa(y) (\nabla_y \omega_i(y) + e_i) \cdot (\nabla_y \omega_j(y) + e_j) dy$.

Only if the product $\phi(\cdot) \tau(\cdot)$ is constant do we get $u_0(t, x, y) = u_0(t, x)$ independent of $y \in Y$, and in that case we can eliminate v_0 from the system to obtain the upscaled pseudoparabolic equation

$$\phi^* \frac{\partial u_0(t, x)}{\partial t} - \nabla \cdot \kappa^* \nabla u_0(t, x) - \nabla \cdot \kappa^* \nabla \phi^* \tau^* \frac{\partial u_0(t, x)}{\partial t} = 0. \quad (9)$$

The homogenized porosity is the average $\phi^* = \int_Y \phi(y) dy$ and the homogenized time-delay is the *harmonic average* $\tau^* = \left(\int_Y \frac{1}{\tau(y)} dy \right)^{-1}$. In the general situation, u_0 depends also on the local variable $y \in Y$, and then the limit system (8) is *partially upscaled*, a combination of the local equations (8a) and the upscaled (8b). We will make similar but much more interesting calculations below when $\phi(\cdot)$ and $\tau(\cdot)$ are *piecewise constant*.

3. The Pseudoparabolic System

Next we extend the models to include binary media of classical or highly-heterogeneous type, and then we obtain the homogenized limit problems by two-scale convergence.

3.1. The Heterogeneous Micro-models

We use a binary medium to emphasize the dependence of singularities on geometry. Let the unit cube Y be given in open disjoint complementary parts, Y_1 and Y_2 , so $Y_1 \cap Y_2 = \emptyset$ and Y is the interior of $\overline{Y_1} \cup \overline{Y_2}$. We denote by $\chi_j(y)$ the characteristic function of Y_j for $j = 1, 2$, extended Y -periodically to all of \mathbb{R}^N . Thus, $\chi_1(y) + \chi_2(y) = 1$ for a.e. y in \mathbb{R}^N . It is assumed that the sets $\{y \in \mathbb{R}^N : \chi_j(y) = 1\}$ for $j = 1, 2$, have smooth boundary, but we do not require these sets to be connected. The corresponding ε -periodic characteristic functions are defined by

$$\chi_j^\varepsilon(x) \equiv \chi_j\left(\frac{x}{\varepsilon}\right), \quad x \in \mathbb{R}^N, \quad j = 1, 2,$$

and these naturally partition the global domain Ω into two sub-domains, Ω_1^ε and Ω_2^ε by $\Omega_j^\varepsilon \equiv \left\{x \in \Omega : \chi_j^\varepsilon(x) = 1\right\}$, $j = 1, 2$. We use the characteristic functions as multipliers to denote the *zero-extension* of various functions. Let $\Gamma \equiv \partial Y_1 \cap \partial Y_2 \cap Y$ be the part of the interface between Y_1 and Y_2 that is interior to the local cell Y . Then $\Gamma^\varepsilon \equiv \partial \Omega_1^\varepsilon \cap \partial \Omega_2^\varepsilon \cap \Omega$ represents the corresponding interface between Ω_1^ε and Ω_2^ε that is interior to Ω . We denote by γ_j the *boundary trace* of functions on Y_j to Γ and by γ_j^ε the *boundary trace* of functions on Ω_j^ε to Γ^ε . (See [33], [29].)

3.1.1. The Classical Case

Let the strictly positive lower-bounded functions $\phi_j(\cdot, \cdot)$, $\kappa_j(\cdot, \cdot)$, $\tau_j(\cdot, \cdot) \in L^\infty(\Omega; C(\overline{Y_j}))$ be given, and define Y -periodic functions in $L^\infty(\Omega; L_{\#}^2(Y))$ by

$$\begin{aligned} \phi(x, y) \equiv \phi_j(x, y), \quad \kappa(x, y) \equiv \kappa_j(x, y), \quad \tau(x, y) \equiv \tau_j(x, y), \\ y \in Y_j, \quad j = 1, 2, \quad x \in \Omega. \end{aligned}$$

The subscript $\#$ denotes the subspace of Y -periodic functions in any function space. Corresponding functions on Ω_j^ε are defined by

$$\begin{aligned} \phi_j^\varepsilon(x) \equiv \phi_j\left(x, \frac{x}{\varepsilon}\right), \quad \kappa_j^\varepsilon(x) \equiv \kappa_j\left(x, \frac{x}{\varepsilon}\right), \quad \tau_j^\varepsilon(x) \equiv \tau_j\left(x, \frac{x}{\varepsilon}\right), \\ x \in \Omega_j^\varepsilon, \quad j = 1, 2, \end{aligned}$$

and the coefficients for the pseudoparabolic system (5) are given by

$$\phi^\varepsilon(x) \equiv \chi_1^\varepsilon(x)\phi_1^\varepsilon(x) + \chi_2^\varepsilon(x)\phi_2^\varepsilon(x), \tag{10a}$$

$$\kappa^\varepsilon(x) \equiv \chi_1^\varepsilon(x)\kappa_1^\varepsilon(x) + \chi_2^\varepsilon(x)\kappa_2^\varepsilon(x), \tag{10b}$$

$$\tau^\varepsilon(x) \equiv \chi_1^\varepsilon(x)\tau_1^\varepsilon(x) + \chi_2^\varepsilon(x)\tau_2^\varepsilon(x). \tag{10c}$$

These are ε -periodic on the *fine scale*. Theorem 1.1 gives a unique solution of the ε -problem: $u^\varepsilon(\cdot) \in H^1((0, T); L^2(\Omega))$ and $v^\varepsilon(\cdot) \in L^2((0, T); H_0^1(\Omega))$ satisfy

$$\int_{\Omega} (\phi^\varepsilon(x) \frac{\partial u^\varepsilon(t, x)}{\partial t} \varphi(x) + \frac{1}{\tau^\varepsilon(x)} (u^\varepsilon(t, x) - v^\varepsilon(t, x)) (\varphi(x) - \psi(x)) + \kappa^\varepsilon(x) \nabla v^\varepsilon(t, x) \cdot \nabla \psi(x)) dx = 0 \tag{11}$$

for all $\varphi(\cdot) \in L^2(\Omega)$ and $\psi(\cdot) \in H_0^1(\Omega)$, together with the initial condition $u^\varepsilon(0, \cdot) = u_*(\cdot)$. The initial value u_* is independent of ε .

If the coefficients κ_j^ε are continuous on Ω_j^ε , the strong form of (11) is the *transmission problem*

$$\phi^\varepsilon(x) \frac{\partial u^\varepsilon(t, x)}{\partial t} + \frac{1}{\tau^\varepsilon(x)} (u^\varepsilon(t, x) - v^\varepsilon(t, x)) = 0, \quad x \in \Omega, \tag{12a}$$

$$-\nabla \cdot (\kappa_1^\varepsilon(x) \nabla v^\varepsilon(t, x)) + \frac{1}{\tau_1^\varepsilon(x)} (v^\varepsilon(t, x) - u^\varepsilon(t, x)) = 0, \quad x \in \Omega_1^\varepsilon, \tag{12b}$$

$$-\nabla \cdot (\kappa_2^\varepsilon(x) \nabla v^\varepsilon(t, x)) + \frac{1}{\tau_2^\varepsilon(x)} (v^\varepsilon(t, x) - u^\varepsilon(t, x)) = 0, \quad x \in \Omega_2^\varepsilon, \tag{12c}$$

$$\gamma_1^\varepsilon v^\varepsilon(t, s) = \gamma_2^\varepsilon v^\varepsilon(t, s), \tag{12d}$$

$$\kappa_1^\varepsilon(s) \nabla v^\varepsilon(t, s) \cdot \nu = \kappa_2^\varepsilon(s) \nabla v^\varepsilon(t, s) \cdot \nu, \quad s \in \Gamma^\varepsilon, \tag{12e}$$

where ν denotes the unit outward normal on $\partial\Omega_1^\varepsilon$. We have homogeneous Dirichlet boundary conditions

$$v^\varepsilon(t, s) = 0 \quad \text{a.e. } s \in \partial\Omega, \tag{12f}$$

and the initial condition $u^\varepsilon(0, x) = u_*(x)$, a.e. $x \in \Omega$. This is the *exact micro-model*. If κ^ε is continuous on Γ^ε , there are no interface conditions and (12) reduces to the single system (5) over Ω . Even then, the fine-scale dependence on the coefficients and geometry make it numerically intractable for realistically small values of $\varepsilon > 0$.

3.1.2. The Highly-Heterogeneous Case

In the highly-heterogeneous case, the permeability is scaled by ε^2 in the second region Ω_2^ε , so the flux is given by $-\varepsilon^2 \kappa_2^\varepsilon(x) \nabla v^\varepsilon$ in Ω_2^ε :

$$\kappa^\varepsilon(x) \equiv \chi_1^\varepsilon(x)\kappa_1^\varepsilon(x) + \varepsilon^2 \chi_2^\varepsilon(x)\kappa_2^\varepsilon(x). \tag{13}$$

Then the system (11) becomes

$$\phi^\varepsilon(x) \frac{\partial u^\varepsilon(t, x)}{\partial t} + \frac{1}{\tau^\varepsilon(x)} (u^\varepsilon(t, x) - v^\varepsilon(t, x)) = 0, \quad x \in \Omega, \quad (14a)$$

$$-\nabla \cdot (\kappa_1^\varepsilon(x) \nabla v^\varepsilon(t, x)) + \frac{1}{\tau_1^\varepsilon(x)} (v^\varepsilon(t, x) - u^\varepsilon(t, x)) = 0, \quad x \in \Omega_1^\varepsilon, \quad (14b)$$

$$-\nabla \cdot (\varepsilon^2 \kappa_2^\varepsilon(x) \nabla v^\varepsilon(t, x)) + \frac{1}{\tau_2^\varepsilon(x)} (v^\varepsilon(t, x) - u^\varepsilon(t, x)) = 0, \quad x \in \Omega_2^\varepsilon, \quad (14c)$$

$$\gamma_1^\varepsilon v^\varepsilon(t, s) = \gamma_2^\varepsilon v^\varepsilon(t, s), \quad (14d)$$

$$\kappa_1^\varepsilon(s) \nabla v^\varepsilon(t, s) \cdot \nu = \varepsilon^2 \kappa_2^\varepsilon(s) \nabla v^\varepsilon(t, s) \cdot \nu, \quad s \in \Gamma^\varepsilon. \quad (14e)$$

The ε -problem for the model developed by Arbogast, Douglas, and Hornung [34] is recovered by letting $\tau^\varepsilon \rightarrow 0$.

3.2. Homogenization of the Classical Case

3.2.1. The Two-scale Limit

Let the coefficients in (5) be given by (10). Denote the gradient in the y -variable by ∇_y , and use the symbol “ $\xrightarrow{2}$ ” to denote *two-scale convergence* [35].

Lemma 3.1: *For each $\varepsilon > 0$, let $u^\varepsilon(\cdot)$, $v^\varepsilon(\cdot)$ denote the unique solution to the pseudoparabolic ε -problem (11). These satisfy the estimates*

$$\|u^\varepsilon\|_{L^2((0,T) \times \Omega)} + \|v^\varepsilon\|_{L^2((0,T); H_0^1(\Omega))} \leq C,$$

so there exist

- (i) a function U in $L^2((0, T) \times \Omega; L^2_\#(Y))$,
- (ii) a function v in $L^2((0, T); H_0^1(\Omega))$,
- (ii) a function V in $L^2((0, T) \times \Omega; H^1_\#(Y)/\mathbb{R})$,

and a subsequence, hereafter denoted by u^ε , v^ε , which two-scale converges as follows:

$$u^\varepsilon \xrightarrow{2} U(t, x, y), \quad (15a)$$

$$v^\varepsilon \xrightarrow{2} v(t, x), \quad (15b)$$

$$\nabla v^\varepsilon \xrightarrow{2} \nabla v(t, x) + \nabla_y V(t, x, y). \quad (15c)$$

This suggests use of the corresponding test functions

$$\tilde{\varphi}(x) = \Phi(x, x/\varepsilon), \quad \tilde{\psi}(x) = \psi(x) + \varepsilon \Psi(x, x/\varepsilon),$$

where $\psi \in H_0^1(\Omega)$, $\Phi, \Psi \in C_0^\infty(\Omega; C_\#^\infty(Y))$. Setting these in (11), we obtain

$$\begin{aligned} & \int_{\Omega} (\phi^\varepsilon(x) \frac{\partial u^\varepsilon(t, x)}{\partial t} \Phi(x, x/\varepsilon) \\ & + \frac{1}{\tau^\varepsilon(x)} (u^\varepsilon(t, x) - v^\varepsilon(t, x)) (\Phi(x, x/\varepsilon) - (\psi(x) + \varepsilon \Psi(x, x/\varepsilon))) \\ & + \kappa^\varepsilon(x) \nabla v^\varepsilon(t, x) \cdot \nabla (\psi(x) + \varepsilon \Psi(x, x/\varepsilon))) dx = 0. \end{aligned}$$

Take the limit as $\varepsilon \rightarrow 0$ to obtain the *two-scale limit system*

$$\begin{aligned} & \int_{\Omega} \int_Y (\phi(x, y) \frac{\partial U(t, x, y)}{\partial t} \Phi(x, y) \\ & + \frac{1}{\tau(x, y)} (U(t, x, y) - v(t, x)) (\Phi(x, y) - \psi(x)) \\ & + \kappa(x, y) (\nabla v(t, x) + \nabla_y V(t, x, y)) \cdot (\nabla \psi(x) + \nabla_y \Psi(x, y))) dy dx = 0. \end{aligned} \quad (16)$$

for all Φ, ψ, Ψ as above, and $U(0, x, y) = u_*(x)$. From the uniqueness of the solution of the initial-value-problem for (16), it follows that the original sequence $u^\varepsilon, v^\varepsilon$ two-scale converges as above.

In order to eliminate the function $V(t, x, y)$ from this system, we use the *periodic cell problem*: for each $k = 1, 2, \dots, N$, define ω_k by

$$\begin{aligned} & \omega_k \in L^2(\Omega; H_\#^1(Y)) : \\ & \int_Y \kappa(x, y) (\nabla_y \omega_k(x, y) + e_k) \cdot \nabla_y \Psi(x, y) dy = 0 \text{ for all } \Psi \in L^2(\Omega; H_\#^1(Y)). \end{aligned} \quad (17)$$

(Let's ask that $\int_Y \omega_k(x, y) dy = 0$ to fix the constant.) Then we have the representation $V(t, x, y) = \sum_{i=1}^N \frac{\partial v(t, x)}{\partial x_i} \omega_i(x, y)$. Specify similar test functions $\Psi(x, y) = \sum_{j=1}^N \frac{\partial \psi(x)}{\partial x_j} \omega_j(x, y)$ above to obtain

Theorem 3.2: *The limits U, v in Lemma 3.1 are the solution of the partially homogenized pseudoparabolic system*

$$\begin{aligned} & U \in H^1((0, T); L^2(\Omega; L_\#^2(Y))), \quad v \in L^2((0, T); H_0^1(\Omega)) : \\ & \int_{\Omega} \int_Y (\phi(x, y) \frac{\partial U(t, x, y)}{\partial t} \Phi(x, y) \\ & + \frac{1}{\tau(x, y)} (U(t, x, y) - v(t, x)) (\Phi(x, y) - \psi(x))) dy dx \\ & + \int_{\Omega} \left(\sum_{i,j=1}^N \kappa_{ij}^*(x) \frac{\partial v(x)}{\partial x_i} \frac{\partial \psi(x)}{\partial x_j} \right) dx = 0, \\ & \text{for all } \Phi \in L^2(\Omega; L_\#^2(Y)), \quad \psi \in H_0^1(\Omega), \end{aligned} \quad (18)$$

and $U(0, x, y) = u_*(x)$, where the effective coefficients are given by

$$\kappa_{ij}^*(x) = \int_Y \kappa(x, y) (\nabla_y \omega_i(x, y) + e_i) \cdot (\nabla_y \omega_j(x, y) + e_j) dy.$$

3.2.2. Summary

The strong formulation of the system (18) is

$$\phi(x, y) \frac{\partial U(t, x, y)}{\partial t} + \frac{1}{\tau(x, y)} (U(t, x, y) - v(t, x)) = 0, \quad y \in Y, \quad (19a)$$

$$\int_Y \frac{1}{\tau(x, y)} (v(t, x) - U(t, x, y)) dy - \nabla \cdot \kappa^* \nabla v(t, x) = 0. \quad (19b)$$

This extends (8) from ε -periodic coefficients to those which depend also on the slow variable, $x \in \Omega$.

Consider the case of a binary medium in which each of $\phi_j, \tau_j \in L^\infty(\Omega)$ is independent of $y \in Y_j$. Then the same is true of

$$U(t, x, y) \equiv \begin{cases} U_1(t, x), & y \in Y_1, \\ U_2(t, x), & y \in Y_2, \end{cases}$$

and we have the homogenized *binary system*

$$|Y_1| \phi_1(x) \frac{\partial U_1(t, x)}{\partial t} + \frac{|Y_1|}{\tau_1(x)} (U_1(t, x) - v(t, x)) = 0, \quad (20a)$$

$$|Y_2| \phi_2(x) \frac{\partial U_2(t, x)}{\partial t} + \frac{|Y_2|}{\tau_2(x)} (U_2(t, x) - v(t, x)) = 0, \quad (20b)$$

$$\begin{aligned} \frac{|Y_1|}{\tau_1(x)} (v(t, x) - U_1(t, x)) + \frac{|Y_2|}{\tau_2(x)} (v(t, x) - U_2(t, x)) \\ - \nabla \cdot \kappa^* \nabla v(t, x) = 0. \end{aligned} \quad (20c)$$

This is the binary medium analogue of (9).

3.3. Homogenization of the Highly-heterogeneous Case

3.3.1. The Two-scale Limit

Here the permeability is given by (13), so we obtain weaker a-priori estimates and correspondingly weaker convergence results.

Lemma 3.3: *For each $\varepsilon > 0$, let $u^\varepsilon(\cdot), v^\varepsilon(\cdot)$ denote the unique solution to the pseudoparabolic ε -problem (11). These satisfy the estimates*

$$\begin{aligned} \|u^\varepsilon\|_{L^2((0,T) \times \Omega)} + \|v^\varepsilon\|_{L^2((0,T) \times \Omega)} \\ + \|v^\varepsilon\|_{L^2((0,T); H^1(\Omega_1^\varepsilon))} + \|\varepsilon v^\varepsilon\|_{L^2((0,T); H^1(\Omega_2^\varepsilon))} \leq C, \end{aligned}$$

so there exist

- (i) a function U in $L^2((0, T) \times \Omega; L^2_{\#}(Y))$,
- (ii) a function v_1 in $L^2((0, T); H^1_0(\Omega))$,
- (iii) a pair of functions V_j in $L^2((0, T) \times \Omega; H^1_{\#}(Y_j)/\mathbb{R})$, $j = 1, 2$,

and a subsequence, hereafter denoted by u^ε , v^ε , which two-scale converges as follows:

$$u^\varepsilon(t, x) \xrightarrow{2} U(t, x, y), \quad (21a)$$

$$\chi_1^\varepsilon v^\varepsilon \xrightarrow{2} \chi_1(y) v_1(t, x), \quad (21b)$$

$$\chi_1^\varepsilon \nabla v^\varepsilon \xrightarrow{2} \chi_1(y) [\nabla v_1(t, x) + \nabla_y V_1(t, x, y)], \quad (21c)$$

$$\chi_2^\varepsilon v^\varepsilon \xrightarrow{2} \chi_2(y) V_2(t, x, y), \quad (21d)$$

$$\varepsilon \chi_2^\varepsilon \nabla v^\varepsilon \xrightarrow{2} \chi_2(y) \nabla_y V_2(t, x, y). \quad (21e)$$

The function V_2 satisfies $\gamma_2(V_2(t, x, y) = v_1(x)$, $y \in \Gamma$. (See [36].) These suggest use of the corresponding test functions

$$\varphi(x) = \Phi(x, x/\varepsilon), \quad \psi(x) = \begin{cases} \psi_1(x) + \varepsilon \Psi_1(x, x/\varepsilon) & : x \in \Omega_1^\varepsilon, \\ \Psi_2(x, x/\varepsilon) + \varepsilon \Psi_1(x, x/\varepsilon) & : x \in \Omega_2^\varepsilon, \end{cases}$$

where $\psi_1 \in H^1_0(\Omega)$, $\Phi, \Psi_1 \in C^\infty_0(\Omega; C^\infty_{\#}(Y))$ and $\Psi_2 \in C^\infty_0(\Omega; C^\infty_{\#}(Y_2))$ with $\gamma_2 \Psi_2(x, \cdot) = \psi_1(x)$ on Γ . Setting these in (11) yields

$$\begin{aligned} & \int_{\Omega} \left(\phi^\varepsilon \frac{\partial u^\varepsilon(t, x)}{\partial t} \Phi(x, x/\varepsilon) \right. \\ & \quad + \frac{\chi_1^\varepsilon(x)}{\tau_1^\varepsilon(x)} (u^\varepsilon(t, x) - v^\varepsilon(t, x)) (\Phi(x, x/\varepsilon) - (\psi_1(x) + \varepsilon \Psi_1(x, x/\varepsilon))) \\ & \quad + \frac{\chi_2^\varepsilon(x)}{\tau_2^\varepsilon(x)} (u^\varepsilon(t, x) - v^\varepsilon(t, x)) (\Phi(x, x/\varepsilon) - (\Psi_2(x, x/\varepsilon) + \varepsilon \Psi_1(x, x/\varepsilon))) \\ & \quad + \chi_1^\varepsilon(x) \kappa_1^\varepsilon(x) \nabla v^\varepsilon(t, x) \cdot \nabla (\psi_1(x) + \varepsilon \Psi_1(x, x/\varepsilon)) \\ & \quad \left. + \chi_2^\varepsilon(x) \kappa_2^\varepsilon(x) \varepsilon \nabla v^\varepsilon(t, x) \cdot \varepsilon \nabla (\Psi_2(x, x/\varepsilon) + \varepsilon \Psi_1(x, x/\varepsilon)) \right) dx = 0. \end{aligned}$$

Take the limit as $\varepsilon \rightarrow 0$ to obtain the *two-scale limit system*

$$\begin{aligned} & \int_{\Omega} \int_Y \left(\phi(x, y) \frac{\partial U(t, x, y)}{\partial t} \Phi(x, y) \right. \\ & \quad + \frac{\chi_1(y)}{\tau_1(x, y)} (U(t, x, y) - v_1(t, x)) (\Phi(x, y) - \psi_1(x)) \\ & \quad + \frac{\chi_2(y)}{\tau_2(x, y)} (U(t, x, y) - V_2(t, x)) (\Phi(x, y) - \Psi_2(x, y)) \\ & \quad + \chi_1(y) \kappa_1(x, y) (\nabla v_1(t, x) + \nabla_y V_1(t, x, y)) \cdot (\nabla \psi_1(x) + \nabla_y \Psi_1(x, y)) \\ & \quad \left. + \chi_2(y) \kappa_2(x, y) \nabla_y V_2(t, x, y) \cdot \nabla_y \Psi_2(x, y) \right) dy dx = 0, \quad (22) \end{aligned}$$

for all Φ , ψ_1 , Ψ_1 , Ψ_2 as above, and $U(0, x, y) = u_*(x)$. The uniqueness of the

solution to the corresponding initial-value problem shows the original sequence converges to it.

As before, we can represent each $V_1(t, x, \cdot)$ by a *cell problem*: define $\omega_k(x, y)$ by

$$\omega_k \in L^2(\Omega; H_{\#}^1(Y_1)) : \int_{Y_1} \kappa_1(x, y) (\nabla_y \omega_k(x, y) + \mathbf{e}_k) \cdot \nabla_y \Psi_1(x, y) dy = 0$$

$$\text{for all } \Psi_1 \in L^2(\Omega; H_{\#}^1(Y_1)), \quad \int_{Y_1} \omega_k(x, y) dy = 0. \quad (23)$$

Then we have $V_1(t, x, y) = \sum_{i=1}^N \frac{\partial v_1(t, x)}{\partial x_i} \omega_i(x, y)$, and we specify the test functions $\Psi_1(x, y) = \sum_{j=1}^N \frac{\partial \psi_1(x)}{\partial x_j} \omega_j(x, y)$ above to obtain

Theorem 3.4: *The limits U, v_1, V_2 in Lemma 3.3 are the solution of the partially homogenized pseudoparabolic system*

$$U \in H^1((0, T); L^2(\Omega; L_{\#}^2(Y))), \quad v_1 \in L^2((0, T); H_0^1(\Omega)),$$

$$V_2 \in L^2((0, T) \times \Omega; H_{\#}^1(Y_2)) \text{ with } \gamma V_2|_{\Gamma} = v_1 :$$

$$\int_{\Omega} \int_Y (\phi(x, y) \frac{\partial U(t, x, y)}{\partial t} \Phi(x, y) + \frac{\chi_1(y)}{\tau_1(x, y)} (U(t, x, y) - v_1(t, x)) (\Phi(x, y) - \psi_1(x))$$

$$+ \frac{\chi_2(y)}{\tau_2(x, y)} (U(t, x, y) - V_2(t, x, y)) (\Phi(x, y) - \Psi_2(x, y))) dy dx$$

$$+ \int_{\Omega} \left(\sum_{i,j=1}^N \kappa_{ij}^*(x) \frac{\partial v_1(t, x)}{\partial x_i} \frac{\partial \psi_1(x)}{\partial x_j} \right) dx + \int_{\Omega} \int_{Y_2} \kappa_2(x, y) \nabla_y V_2(t, x, y) \cdot \nabla_y \Psi_2(x, y) dy dx = 0,$$

$$\text{for all } \Phi \in L^2(\Omega; L_{\#}^2(Y)), \quad \psi_1 \in H_0^1(\Omega),$$

$$\Psi_2 \in L^2(\Omega; H_{\#}^1(Y_2)) \text{ with } \gamma \Psi_2|_{\Gamma} = \psi_1, \quad (24)$$

and $U(0, x, y) = u_*(x)$, where the effective coefficients are given by

$$\kappa_{ij}^*(x) = \int_{Y_1} \kappa_1(x, y) (\nabla_y \omega_i(x, y) + \mathbf{e}_i) \cdot (\nabla_y \omega_j(x, y) + \mathbf{e}_j) dy.$$

Next we separate the components of the system. First write the part over Y_2

$$\phi_2(x, y) \frac{\partial U(t, x, y)}{\partial t} + \frac{1}{\tau_2(x, y)} (U(t, x, y) - V_2(t, x, y)) = 0 \text{ and}$$

$$\frac{1}{\tau_2(x, y)} (V_2(t, x, y) - U(t, x, y)) - \nabla_y \cdot \kappa_2(x, y) \nabla_y V_2(t, x, y) = 0, \quad y \in Y_2,$$

$$\gamma V_2(t, x, y) = v_1(t, x), \quad y \in \Gamma,$$

and then substitute these back into (24) and use Stokes' theorem on Y_2 to get

$$\begin{aligned} & \int_{\Omega} \int_{Y_1} (\phi_1(x, y) \frac{\partial U(t, x, y)}{\partial t} \Phi(x, y) \\ & \quad + \frac{1}{\tau_1(x, y)} (U(t, x, y) - v_1(t, x)) (\Phi(x, y) - \psi_1(x))) dy dx \\ & + \int_{\Omega} \left(\sum_{i,j=1}^N \kappa_{ij}^*(x) \frac{\partial v_1(t, x)}{\partial x_i} \frac{\partial \psi_1(x)}{\partial x_j} \right) dx + \int_{\Omega} \int_{\Gamma} \kappa_2(x, y) \nabla_y V_2(t, x, y) \cdot \nu dS \psi_1(x) dx = 0. \end{aligned}$$

3.3.2. Summary

The strong form of the partially homogenized system (24) is

$$\begin{aligned} \phi_1(x, y) \frac{\partial U(t, x, y)}{\partial t} + \frac{1}{\tau_1(x, y)} (U(t, x, y) - v_1(t, x)) &= 0, \quad y \in Y_1, \\ \int_{Y_1} \frac{1}{\tau_1(x, y)} (v_1(t, x) - U(t, x, y)) dy - \nabla \cdot \kappa^* \nabla v_1(t, x) & \quad (25a) \\ + \int_{\Gamma} \kappa_2(x, y) \nabla_y V_2(t, x, y) \cdot \nu dS &= 0, \end{aligned}$$

and for each $x \in \Omega$,

$$\begin{aligned} \phi_2(x, y) \frac{\partial U(t, x, y)}{\partial t} + \frac{1}{\tau_2(x, y)} (U(t, x, y) - V_2(t, x, y)) &= 0, \\ \frac{1}{\tau_2(x, y)} (V_2(t, x, y) - U(t, x, y)) - \nabla_y \cdot \kappa_2(x, y) \nabla_y V_2(t, x, y) &= 0, \quad y \in Y_2, \\ \gamma V_2(t, x, y) = v_1(t, x), \quad y \in \Gamma. & \quad (25b) \end{aligned}$$

Note the coupling in the system: the function v_1 from (25a) is input to (25b), and the total flux from (25b) is the distributed source in (25a).

Suppose now that ϕ_1 and τ_1 are independent of $y \in Y_1$, and therefore so also is $u(t, x) \equiv U(t, x, y)$, $y \in Y_1$. Then (25a) is homogenized:

$$\begin{aligned} \phi_1(x) \frac{\partial u(t, x)}{\partial t} + \frac{1}{\tau_1(x)} (u(t, x) - v_1(t, x)) &= 0, \\ \frac{1}{\tau_1(x)} (v_1(t, x) - u(t, x)) - \frac{1}{|Y_1|} \nabla \cdot \kappa^* \nabla v_1(t, x) & \quad (26a) \\ + \frac{1}{|Y_1|} \int_{\Gamma} \kappa_2(x, y) \nabla_y V_2(t, x, y) \cdot \nu dS &= 0, \end{aligned}$$

and for each $x \in \Omega$,

$$\begin{aligned} \phi_2(x, y) \frac{\partial U(t, x, y)}{\partial t} + \frac{1}{\tau_2(x, y)} (U(t, x, y) - V_2(t, x, y)) &= 0, \\ \frac{1}{\tau_2(x, y)} (V_2(t, x, y) - U(t, x, y)) - \nabla_y \cdot \kappa_2(x, y) \nabla_y V_2(t, x, y) &= 0, \quad y \in Y_2, \\ \gamma V_2(t, x, y) &= v_1(t, x), \quad y \in \Gamma. \end{aligned} \quad (26b)$$

Note that (26a) is the upscaled fissured medium system, and (26b) is the local fissured medium system at each $x \in \Omega$.

3.4. Vanishing time delay

Suppose that $\tau_1^\varepsilon = o(\varepsilon)$ in the classical system (12). Then $\|u^\varepsilon - v^\varepsilon\|_{L^2(Y_1)} = o(\varepsilon^{1/2})$, so in the limit we obtain $U(t, x, y)|_{Y_1} = v(t, x)$. Choose test functions $\Phi(x, y) = \psi(x) + \varepsilon \Psi(x, y)$ in the weak form, with the equations added, and take the limit to get the homogenized mixed parabolic-pseudoparabolic system (compare (20))

$$\phi_1^*(x) \frac{\partial v(t, x)}{\partial t} - \nabla \cdot \kappa^* \nabla v(t, x) + \int_{Y_2} \frac{1}{\tau_2(x, y)} (v(t, x) - U(t, x, y)) dy = 0, \quad (27a)$$

$$\phi_2(x, y) \frac{\partial U(t, x, y)}{\partial t} + \frac{1}{\tau_2(x, y)} (U(t, x, y) - v(t, x)) = 0, \quad y \in Y_2, \quad (27b)$$

with effective porosity $\phi_1^*(x) = \int_{Y_1} \phi_1(x, y) dy$. Then (27a) is a parabolic equation with a memory term determined by (27b). See Peszyńska [37] for results and additional references to memory functionals in parabolic equations; also see [31] for first-order kinetic models.

Suppose that $\tau_1^\varepsilon = o(\varepsilon)$ in the highly-heterogeneous system (14). Then $U(t, x, y)|_{Y_1} = v_1(t, x)$ and instead of the system (25a) we obtain the homogenized parabolic equation

$$\phi_1^*(x) \frac{\partial v_1(t, x)}{\partial t} - \nabla \cdot \kappa^* \nabla v_1(t, x) + \int_{\Gamma} \kappa_2(x, y) \nabla_y V_2(x, y) \cdot \nu dS = 0. \quad (28a)$$

Suppose that $\tau_2^\varepsilon = o(\varepsilon)$ in (14). Then $U(t, x, y)|_{Y_2} = V_2(t, x, y)$ and instead of the system (25b) we obtain the local parabolic equations

$$\phi_2(x, y) \frac{\partial V_2(t, x, y)}{\partial t} - \nabla_y \cdot \kappa_2(x, y) \nabla_y V_2(t, x, y) = 0, \quad y \in Y_2, \quad (28b)$$

$$\gamma V_2(x, y) = v_1(x), \quad y \in \Gamma. \quad (28c)$$

If both vanish in the limit, then we recover the Arbogast-Douglas-Hornung [34] double-porosity model (28) of a fractured porous medium.

4. Partially-saturated Flow with Dynamic Capillary Pressure

4.1. Microscopic Equations

Let us consider the unsaturated flow in a highly-heterogeneous medium Ω with the ε -periodic structure of Section 3. Here Y_2 is the matrix block and Y_1 is the surrounding fracture domain. Each of the subdomains Ω_i^ε is characterized by a rock permeability tensor K^i , a porosity ϕ^i , the relative permeability $k_w^i(u^i)$, and the capillary pressure function $P_c^i(u^i)$. Here u^i denotes the saturation in Ω_i^ε . The fluid has constant viscosity μ and density ρ . It has been observed that the dynamic effects in capillary pressure equilibrium are much more significant in media with low conductivity than those with high conductivity, so we assume that the unsaturated flow can be locally described by the original Richards equation (1) in the fracture domain Ω_1^ε and by the pseudoparabolic Richards equation (2) in the porous matrix Ω_2^ε :

$$\phi^1 \frac{\partial u^1}{\partial t} + \nabla \cdot \frac{1}{\mu} K^1 k_w^1(u^1) \nabla (P_c^1(u^1) - \rho G D(x)) = 0, \quad x \in \Omega_1^\varepsilon, \quad (29a)$$

$$\begin{aligned} \phi^2 \frac{\partial u^2}{\partial t} + \varepsilon^2 \nabla \cdot \frac{1}{\mu} K^2 k_w^2(u^2) \nabla (P_c^2(u^2) - \tau \frac{\partial u^2}{\partial t} - \rho G D(x)) \\ = 0, \quad x \in \Omega_2^\varepsilon. \end{aligned} \quad (29b)$$

Hereafter for simplicity we set depth $D(x) = x_3$. Introduce $p^i = -P_c^i(u^i)$, $u^i = \alpha^i(p^i)$, $\kappa^i(p^i) = \frac{1}{\mu} K^i k_w^i(u^i)$, so $\alpha^i(\cdot)$ is inverse to $-P_c^i(\cdot)$, and equations (29a) and (29b) can be rewritten as

$$\phi^1 \frac{\partial \alpha^1(p^1)}{\partial t} - \nabla \cdot \kappa^1(p^1) (\nabla p^1 + \rho G \mathbf{e}_3) = 0, \quad x \in \Omega_1^\varepsilon, \quad (30a)$$

$$\phi^2 \frac{\partial \alpha^2(p^2)}{\partial t} - \varepsilon^2 \nabla \cdot \kappa^2(p^2) (\nabla p^2 + \tau \nabla \frac{\partial \alpha^2(p^2)}{\partial t} + \rho G \mathbf{e}_3) = 0, \quad x \in \Omega_2^\varepsilon, \quad (30b)$$

and are subject to the interface conditions

$$p^1 = p^2 + \tau \frac{\partial \alpha^2(p^2)}{\partial t}, \quad x \in \Gamma^\varepsilon, \quad (30c)$$

$$\begin{aligned} \kappa^1(p^1) (\nabla p^1 + \rho G \mathbf{e}_3) \cdot \nu \\ = \varepsilon^2 \kappa^2(p^2) (\nabla p^2 + \tau \nabla \frac{\partial \alpha^2(p^2)}{\partial t} + \rho G \mathbf{e}_3) \cdot \nu, \quad x \in \Gamma^\varepsilon, \end{aligned} \quad (30d)$$

where ν is the unit normal on Γ^ε out of Ω_2^ε , and the initial conditions are

$$p^i(x, 0) = p_*^i(x), \quad x \in \Omega_i^\varepsilon, \quad i = 1, 2. \quad (30e)$$

4.2. Asymptotic Expansions

We shall expand the solution in powers of ε in the form

$$p^i(t, x) = p_0^i(t, x, y) + \varepsilon p_1^i(t, x, y) + \varepsilon^2 p_2^i(t, x, y) + \dots, \quad i = 1, 2, \quad (31)$$

where p_k^i are Y -periodic in $y \in Y_i$ for $k = 0, 1, 2, \dots$. Following methods of [38, 39], we develop various nonlinear quantities $\theta(p)$ in powers of ε by

$$\begin{aligned} \theta(p^i) &= \theta(p_0^i) + \theta'(p_0^i)(p^i - p_0^i) + \theta''(p_0^i)(p^i - p_0^i)^2/2 + \dots \\ &= \theta(p_0^i) + \theta'(p_0^i)(\varepsilon p_1^i + \varepsilon^2 p_2^i + \dots) + \theta''(p_0^i)(\varepsilon p_1^i + \varepsilon^2 p_2^i + \dots)^2/2 + \dots \\ &= \theta(p_0^i) + \varepsilon \theta'(p_0^i) p_1^i + \varepsilon^2 (\theta'(p_0^i) p_2^i + \theta''(p_0^i) (p_1^i)^2/2) + \dots \\ &= \theta(p_0^i) + \varepsilon \hat{\theta}_1^i + \varepsilon^2 \hat{\theta}_2^i + \dots, \text{ for appropriate } \hat{\theta}_1^i, \hat{\theta}_2^i, \dots, i = 1, 2. \end{aligned}$$

Now, we substitute (31) into the microscopic model and expand the gradient according to the relation $\nabla = \nabla_x + \frac{1}{\varepsilon} \nabla_y$. Then, we collect terms by powers of ε . From (30a) we obtain three equations for the combined $\varepsilon^{-2}, \varepsilon^{-1}$, and ε^0 terms when $x \in \Omega, y \in Y_1$:

$$\nabla_y \cdot (\kappa^1(p_0^1) \nabla_y p_0^1) = 0, \tag{32a}$$

$$\begin{aligned} \nabla_y \cdot (\kappa^1(p_0^1) (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3) + \hat{\kappa}_1^1 \nabla_y p_0^1) \\ + \nabla_x \cdot (\kappa^1(p_0^1) \nabla_y p_0^1) = 0, \end{aligned} \tag{32b}$$

$$\begin{aligned} \phi^1 \frac{\partial \alpha^1(p_0^1)}{\partial t} - \nabla_x \cdot (\kappa^1(p_0^1) (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3) + \hat{\kappa}_1^1 \nabla_y p_0^1) \\ - \nabla_y \cdot (\kappa^1(p_0^1) (\nabla_x p_1^1 + \nabla_y p_2^1) \\ + \hat{\kappa}_1^1 (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3) + \hat{\kappa}_2^1 \nabla_y p_0^1) = 0. \end{aligned} \tag{32c}$$

First, equations for ε^0 from (30b) and (30c) are for $x \in \Omega$

$$\phi^2 \frac{\partial \alpha^2(p_0^2)}{\partial t} - \nabla_y \cdot \kappa^2(p_0^2) \nabla_y \left(p_0^2 + \tau \frac{\partial \alpha^2(p_0^2)}{\partial t} \right) = 0, \quad y \in Y_2, \tag{33a}$$

$$p_0^2 + \tau \frac{\partial \alpha^2(p_0^2)}{\partial t} = p_0^1, \quad y \in \Gamma. \tag{33b}$$

The $\varepsilon^{-1}, \varepsilon^0$ and ε^1 equations of (30d) for $x \in \Omega, y \in \Gamma$ are

$$\kappa^1(p_0^1) \nabla_y p_0^1 \cdot \nu = 0, \tag{34a}$$

$$(\kappa^1(p_0^1) (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3) + \hat{\kappa}_1^1 \nabla_y p_0^1) \cdot \nu = 0, \tag{34b}$$

$$\begin{aligned} (\kappa^1(p_0^1) (\nabla_x p_1^1 + \nabla_y p_2^1) + \hat{\kappa}_1^1 (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3) + \hat{\kappa}_2^1 \nabla_y p_0^1) \cdot \nu \\ = \kappa^2(p_0^2) \nabla_y \left(p_0^2 + \tau \frac{\partial \alpha^2(p_0^2)}{\partial t} \right) \cdot \nu. \end{aligned} \tag{34c}$$

Equations (32a) and (34a) form an elliptic system for p_0^1 in terms of y . Since its solution is independent of y , it follows that $p_0^1 = p_0^1(t, x)$, so all terms with $\nabla_y p_0^1$ vanish.

The equations (32b) and (34b) form a linear elliptic system in y whose solution p_1^1 can be represented in terms of p_0^1 . Define $\omega_j(y)$ for $j = 1, 2, 3$ as the Y -periodic solution of the cell problem (compare (23))

$$\nabla_y^2 \omega_j = 0 \text{ for } y \in Y_1, \tag{35a}$$

$$\nabla_y \omega_j \cdot \nu = -\mathbf{e}_j \cdot \nu = -\nu_j \text{ for } y \in \Gamma. \tag{35b}$$

Then from equation (32b) we obtain the representation

$$p_1^1(x, y, t) = \sum_{j=1}^3 \omega_j(y) \left(\frac{\partial p_0^1}{\partial x_j}(x, t) + \rho G \delta_{3j} \right). \quad (36)$$

Now, we locally average (32c) by integrating it over Y_1 to remove the y -variable and get

$$\begin{aligned} & |Y_1| \phi^1 \frac{\partial \alpha^1(p_0^1)}{\partial t} - \int_{Y_1} \nabla_x \cdot \kappa^1(p_0^1) (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3) dy \\ &= \int_{Y_1} \nabla_y \cdot (\kappa^1(p_0^1) (\nabla_x p_1^1 + \nabla_y p_2^1) + \hat{\kappa}_1^1 (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3)) dy. \end{aligned} \quad (37)$$

Apply the divergence theorem to the second integral above, use (34c), make a second application of the divergence theorem, and use (33a) to obtain

$$\begin{aligned} & \int_{Y_1} \nabla_y \cdot (\kappa^1(p_0^1) (\nabla_x p_1^1 + \nabla_y p_2^1) + \hat{\kappa}_1^1 (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3)) dy \\ &= \int_{\partial Y_1} (\kappa^1(p_0^1) (\nabla_x p_1^1 + \nabla_y p_2^1) + \hat{\kappa}_1^1 (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3)) \cdot \nu dS \\ &= - \int_{\partial Y_2} \kappa^2(p_0^2) \nabla_y \left(p_0^2 + \tau \frac{\partial \alpha^2(p_0^2)}{\partial t} \right) \cdot \nu dS \\ &= - \int_{Y_2} \nabla_y \cdot \kappa^2(p_0^2) \nabla_y \left(p_0^2 + \tau \frac{\partial \alpha^2(p_0^2)}{\partial t} \right) dy \\ &= - \int_{Y_2} \phi^2 \frac{\partial \alpha^2(p_0^2)}{\partial t} dy. \end{aligned}$$

The first integral in (37) is evaluated using (36). Its integrand becomes (with implied summation)

$$\begin{aligned} & \nabla_x \cdot \kappa^1(p_0^1) (\nabla_x p_0^1 + \nabla_y p_1^1 + \rho G \mathbf{e}_3) \\ &= \frac{\partial}{\partial x_k} \left[\kappa^1(p_0^1) \left(\frac{\partial p_0^1}{\partial x_k} + \frac{\partial \omega_j}{\partial y_k} \left(\frac{\partial p_0^1}{\partial x_j} + \rho G \delta_{3j} \right) + \rho G \delta_{3k} \right) \right] \\ &= \frac{\partial}{\partial x_k} \left[\kappa^1(p_0^1) \left(\frac{\partial \omega_j}{\partial y_k} + \delta_{jk} \right) \left(\frac{\partial p_0^1}{\partial x_j} + \rho G \delta_{3j} \right) \right]. \end{aligned}$$

Define the effective fracture permeability tensor $K^* = \{K_{jk}^*\}$ and the macroscopic fracture porosity ϕ^* by

$$K_{jk}^* = K^1 \int_{Y_1} \left(\frac{\partial \omega_j}{\partial y_k} + \delta_{jk} \right) dy, \quad \phi^* = |Y_1| \phi^1.$$

We also define

$$\kappa^*(p) = \frac{1}{\mu} K^* k_w^1(\alpha^1(p)).$$

Then, the equation for p_0^1 is

$$\phi^* \frac{\partial \alpha^1(p_0^1)}{\partial t} - \nabla_x \cdot \kappa^*(p_0^1)(\nabla_x p_0^1 + \rho G \mathbf{e}_3) = - \int_{Y_2} \phi^2 \frac{\partial \alpha^2(p_0^2)}{\partial t} dy.$$

4.3. Summary

The complete system of flow equations for $p_0^1(x, t)$, $p_0^2(x, y, t)$ is given by

$$\phi^* \frac{\partial \alpha^1(p_0^1)}{\partial t} + \int_{Y_2} \phi^2 \frac{\partial \alpha^2(p_0^2)}{\partial t} dy \tag{38a}$$

$$-\nabla_x \cdot \kappa^*(p_0^1)(\nabla_x p_0^1 + \rho G \mathbf{e}_3) = 0, \quad x \in \Omega,$$

$$\phi^2 \frac{\partial \alpha^2(p_0^2)}{\partial t} - \nabla_y \cdot \kappa^2(p_0^2) \nabla_y \left(p_0^2 + \tau \frac{\partial \alpha^2(p_0^2)}{\partial t} \right) = 0, \quad y \in Y_2, \tag{38b}$$

$$p_0^2 + \tau \frac{\partial \alpha^2(p_0^2)}{\partial t} = p_0^1, \quad y \in \Gamma, \tag{38c}$$

$$p_0^1(x, 0) = p_{\text{init}}^1(x), \quad p_0^2(x, y, 0) = p_{\text{init}}^2(x), \quad y \in Y_2. \tag{38d}$$

This is the double-porosity model consisting of the upscaled equation (38a) together with the distributed family of local boundary-value problems (38b), (38c) for $x \in \Omega$. It is a nonlinear analogue of the system (28a), (26b).

References

- [1] E. Milne, *The diffusion of imprisoned radiation through a gas*, J. London Math. Soc. 1 (1926), pp. 40–51.
- [2] G.I. Barenblatt, I.P. Zheltov, and I.N. Kochina, *Basic concepts in the theory of seepage of homogeneous liquids in fissured rocks (strata)*, J. Appl. Math. Mech. 24 (1960), pp. 1286–1303.
- [3] L.I. Rubenstein, *On the problem of the process of propagation of heat in heterogeneous media*, Izv. Akad. Nauk SSSR, Ser. Geogr. 1 (1948).
- [4] R.E. Showalter, *The final value problem for evolution equations*, J. Math. Anal. Appl. 47 (1974), pp. 563–572.
- [5] M. Burger, G. Gilboa, S. Osher, and J. Xu, *Nonlinear inverse scale space methods*, Commun. Math. Sci. 4 (2006), pp. 179–212.
- [6] B.C. Aslan, W.W. Hager, and S. Moskow, *A generalized eigenproblem for the Laplacian which arises in lightning*, J. Math. Anal. Appl. 341 (2008), pp. 1028–1041.
- [7] R.E. Showalter and T.W. Ting, *Pseudoparabolic partial differential equations*, SIAM J. Math. Anal. 1 (1970), pp. 1–26.
- [8] R.E. Showalter, *Partial differential equations of Sobolev-Galpern type*, Pacific J. Math. 31 (1969), pp. 787–793.
- [9] T.W. Ting, *Parabolic and pseudo-parabolic partial differential equations*, J. Math. Soc. Japan 21 (1969), pp. 440–453.
- [10] R.W. Carroll and R.E. Showalter *Singular and degenerate Cauchy problems*, Academic Press [Harcourt Brace Jovanovich Publishers], New York, 1976 Mathematics in Science and Engineering, Vol. 127.
- [11] R. Collins *Flow of Fluids Through Porous Materials*, Petroleum Publishing Company, Tulsa, 1976 (Originally published by Van Nostrand Reinhold, 1961).
- [12] J.S. Selker, C.K. Keller, and J.T. McCord *Vadose zone processes*, CRC Press LLC, Boca Raton, 1999.
- [13] H.W. Alt and E. DiBenedetto, *Nonsteady flow of water and oil through inhomogeneous porous media*, Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4) 12 (1985), pp. 335–392.
- [14] H. Alt, S. Luckhaus, and A. Visintin, *On nonstationary flow through porous media*, Ann. Mat. Pura Appl. 136 (1984), pp. 303–316.
- [15] T. Arbogast, *The existence of weak solutions to single porosity and simple dual-porosity models of two-phase incompressible flow*, Nonlinear Anal. 19 (1992), pp. 1009–1031.
- [16] S. Hassanizadeh and W.G. Gray, *Thermodynamic basis of capillary pressure on porous media*, Water Resour. Res. 29 (1993), pp. 3389–3405.
- [17] G. Barenblatt and A.A. Gil'man, *A mathematical model of non-equilibrium counter-current capillary imbibition*, J. Eng. Phys. 52 (1987), pp. 456–461.
- [18] G.I. Barenblatt, D.B. Silin, and T.W. Patzek, *The Mathematical Model of Non-Equilibrium Effects in Water-Oil Displacement*, SPEJ 8 (2003), pp. 409–416.

- [19] S. Hassanizadeh, M. Celia, and H. Dahle, *Dynamic effects in the capillary pressure-saturation relationship and their impacts on unsaturated flow*, Vadose Zone Journal 1 (2002), pp. 38–57.
- [20] D. Wildenschild and K. Jensen, *Laboratory investigations of effective flow behavior in unsaturated heterogeneous sands*, Water Resources Research 35 (1999), pp. 17–27.
- [21] D. Wildenschild, J. Hopmans, A. Kent, and M. Rivers, *Quantitative Analysis of Flow Processes in a Sand Using Synchrotron-Based X-ray Microtomography*, Vadose Zone Journal 4 (2005), pp. 112–126.
- [22] D. Wildenschild, J.W. Hopmanns, and J. Simunek, *Flow rate dependence of soil hydraulic characteristics*, Soil Sci. Soc. Am. J. 65 (2001), pp. 35–48.
- [23] A. Bourgeat and M. Panfilov, *Effective two-phase flow through highly heterogeneous porous media: capillary nonequilibrium effects*, Comput. Geosci. 2 (1998), pp. 191–215.
- [24] M. Peszyńska and S.Y. Yi, *Numerical methods for unsaturated flow with dynamic capillary pressure in heterogeneous porous media*, Int. J. Numer. Anal. Model. 5 (2008), pp. 126–149.
- [25] S.M. Hassanizadeh, M. Celia, and H. Dahle, *Dynamic Effect in the Capillary Pressure-Saturation Relationship and its Impacts on Unsaturated Flow*, Vadose Zone Journal 1 (2002), pp. 38–57.
- [26] M. Böhm and R.E. Showalter, *Diffusion in fissured media*, SIAM J. Math. Anal. 16 (1985), pp. 500–509.
- [27] R.E. Showalter, *Degenerate evolution equations and applications*, Indiana Univ. Math. J. 23 (1973/74), pp. 655–677.
- [28] R.E. Showalter *Hilbert space methods for partial differential equations*, Pitman, London, 1977 Monographs and Studies in Mathematics, Vol. 1.
- [29] R.E. Showalter *Monotone operators in Banach space and nonlinear partial differential equations*, Mathematical Surveys and Monographs Vol. 49, American Mathematical Society, Providence, RI, 1997.
- [30] A. Bensoussan, J.L. Lions, and G. Papanicolaou *Asymptotic analysis for periodic structures*, Studies in Mathematics and its Applications Vol. 5, North-Holland Publishing Co., Amsterdam, 1978.
- [31] U. Hornung (ed.) *Homogenization and porous media*, Interdisciplinary Applied Mathematics Vol. 6, Springer-Verlag, New York, 1997.
- [32] V.V. Jikov, S.M. Kozlov, and O.A. Oleinik *Homogenization of differential operators and integral functionals*, Springer-Verlag, Berlin, 1994.
- [33] R.A. Adams *Sobolev spaces*, Academic Press, New York-London, 1975 Pure and Applied Mathematics, Vol. 65.
- [34] T. Arbogast, J. Douglas Jr., and U. Hornung, *Derivation of the double porosity model of single phase flow via homogenization theory*, SIAM J. Math. Anal. 21 (1990), pp. 823–836.
- [35] G. Allaire, *Homogenization and two-scale convergence*, SIAM J. Math. Anal. 23 (1992), pp. 1482–1518.
- [36] G.W. Clark and R.E. Showalter, *Two-scale convergence of a model for flow in a partially fissured medium*, Electron. J. Differential Equations (1999), pp. No. 2, 20 pp. (electronic).
- [37] M. Peszyńska, *On a model of nonisothermal flow through fissured media*, Differential Integral Equations 8 (1995), pp. 1497–1516.
- [38] J. Douglas Jr. and T. Arbogast, *Dual Porosity Models for Flow in Naturally Fractured Reservoirs*, , in *Dynamics of Fluids in Hierarchical Porous Media* Academic Press, London, 1990, chap. VII, pp. 177–220.
- [39] C.J. van Duijn, H. Eichel, R. Helmig, and I.S. Pop, *Effective equations for two-phase flow in porous media: the effect of trapping on the microscale*, Transp. Porous Media 69 (2007), pp. 411–428.