

## REMARKS ON CAUCHY PROBLEM

Let  $A$  denote a linear operator on a Hilbert space  $H$ , with domain  $D(L) \subset H$  a subspace. (Note that we do *not* make any continuity assumptions on  $A$ .) Suppose the operator satisfies

- $(Av, v)_H \geq 0$  for all  $v \in D(L)$ ,
- $I + A$  maps  $D(L)$  onto  $H$ .

If the operator satisfies a  $V$ -elliptic type assumption, then the first condition will follow easily, and then the Lax-Milgram theorem can be used to show that the second condition also holds. Moreover, these conditions hold for much more general classes of operators than those arising from *elliptic* problems.

We will show later that if an operator satisfies these two conditions, then there is a unique function  $u : [0, \infty) \rightarrow H$  which satisfies the *initial-value problem*

$$cu'(t) + Au(t) = 0, \quad u(0) = u_0.$$

Here we assume  $c > 0$  and  $u_0 \in H$  are given. Also, in this situation it will follow that the corresponding non-homogeneous equation is likewise solvable.

EXAMPLE Let  $H = L^2(0, 1)$ ,  $D(L) = \{v \in H^1(0, 1) : v(0) = 0\}$ , and  $A = \frac{d}{dx}$ .

The two conditions are a bit restrictive, but we can relax them considerably with an elementary observation. Suppose that  $u$  is a solution of the initial-value problem above, and let  $\lambda \in \mathbb{R}$ . Define  $w(t) = e^{-\lambda t}u(t)$  for  $t \geq 0$ . Then it is easy to check that  $w(t)$  is a solution of the problem

$$cw'(t) + (\lambda cI + A)w(t) = 0, \quad w(0) = u_0,$$

and, conversely,  $w(t)$  is a solution of this problem only if  $u(t) = e^{\lambda t}w(t)$  is a solution of the original problem. Thus, the initial-value problem is well-posed if

- $((\lambda cI + A)v, v)_H \geq 0$  for all  $v \in D(L)$ ,
- $(1 + \lambda c)I + A$  maps  $D(L)$  onto  $H$ ,

for *some*  $\lambda \in \mathbb{R}$ . Frequently this is satisfied for large enough  $\lambda$ .

Suppose that  $u(t)$  is the solution under these more general hypotheses. Then we find that

$$\frac{d}{dt}c\|u(t)\|_H^2 = -2(Au(t), u(t))_H \leq 2\lambda c\|u(t)\|_H^2,$$

so we conclude that

$$c\|u(t)\|_H^2 \leq e^{2\lambda t}c\|u_0\|_H^2, \quad t \geq 0.$$

This *stability estimate* shows it is worthwhile to know how small the number  $\lambda$  can be taken.

## ADVECTION-DIFFUSION EQUATION

The conservation equation and flux constitutive equation are

$$c\dot{u} + \nabla \cdot \mathbf{j} = F(x), \quad \mathbf{j} = -k\nabla u + \mathbf{b}u.$$

where  $c = c(x)$ ,  $k = k(x)$  and  $\mathbf{b} = \mathbf{b}(x)$ .

**Gravity-driven Fluid Flow.** Assume  $\lambda \geq 0$ . Let  $V$  be a subspace of  $H^1(G)$  that contains  $H_0^1(G)$ . and consider the boundary-value problem

$$u \in V : \int_G (\lambda cuv + (k\nabla u - u\mathbf{b}) \cdot \nabla v) dx = \int_G Fv dx, \quad v \in V.$$

Define the corresponding operator  $\mathcal{A}_\lambda^R : V \rightarrow V'$  by

$$\begin{aligned} \mathcal{A}_\lambda^R u(v) &= \int_G (\lambda cuv + k\nabla u \cdot \nabla v - u\mathbf{b} \cdot \nabla v) dx \\ &= \int_G (\lambda cu - \nabla \cdot (k\nabla u - \mathbf{b}u)) v dx + \int_{\partial G} (k\nabla u - \mathbf{b}u) \cdot \mathbf{n} v dS \end{aligned}$$

This displays the PDE and the boundary conditions.

Estimates:

$$\begin{aligned} \mathcal{A}_\lambda^R u(u) &= \int_G (\lambda cu^2 + k|\nabla u|^2 + \frac{1}{2}\nabla \cdot \mathbf{b}u^2) dx \\ &\quad - \frac{1}{2} \int_{\partial G} u^2 \mathbf{b} \cdot \mathbf{n} dS \end{aligned}$$

$$\mathcal{A}_\lambda^R u(u) \geq 0 \text{ if } \nabla \cdot \mathbf{b} \geq 0 \text{ in } G \text{ and } v|_{\partial G} = 0 \text{ where } \mathbf{b} \cdot \mathbf{n} > 0 \forall v \in V$$

This is the *outflow region*. This problem is appropriate for models of *fluid flow* where  $\mathbf{b}$  is the gravity term.

**Transport of Concentration.**

$$u \in V : \int_G (\lambda cuv + k\nabla u \cdot \nabla v + \nabla \cdot (u\mathbf{b})v) dx = \int_G Fv dx, \quad v \in V.$$

Define the operator  $\mathcal{A}_\lambda^N : V \rightarrow V'$  by

$$\begin{aligned} \mathcal{A}_\lambda^N u(v) &= \int_G (\lambda cuv + k\nabla u \cdot \nabla v + \nabla \cdot (u\mathbf{b})v) dx \\ &= \int_G (\lambda cu - \nabla \cdot (k\nabla u - \mathbf{b}u)) v dx + \int_{\partial G} k\nabla u \cdot \mathbf{n} v dS \end{aligned}$$

This leads to the same PDE but with a modified boundary condition: there is no advection in the boundary condition.

Estimates:

$$\begin{aligned} \mathcal{A}_\lambda^N u(u) &= \int_G (\lambda cu^2 + k|\nabla u|^2 + \frac{1}{2}\nabla \cdot \mathbf{b}u^2) dx \\ &\quad + \frac{1}{2} \int_{\partial G} u^2 \mathbf{b} \cdot \mathbf{n} dS \end{aligned}$$

$$\mathcal{A}_\lambda^N u(u) \geq 0 \text{ if } \nabla \cdot \mathbf{b} \geq 0 \text{ in } G \text{ and } v|_{\partial G} = 0 \text{ where } \mathbf{b} \cdot \mathbf{n} < 0 \forall v \in V$$

This is the *inflow region*. Such problems arise as models of *concentration transport* where  $\mathbf{b}$  is the velocity term. Note that

$$\mathcal{A}_\lambda^N u(v) = \mathcal{A}_\lambda^R u(v) + \int_{\partial G} \mathbf{b} \cdot \mathbf{n} uv dS$$

so it is easier for  $\mathcal{A}_\lambda^N$  to be coercive than for  $\mathcal{A}_\lambda^R$ .